

Technical Paper

HYDRAULICS OF RIVER FLOW UNDER ARCH BRIDGES

-A PROGRESS REPORT-

TO: K. B. Woods, Director

May 14, 1959

Joint Highway Research Project

FROM:

H. L. Michael, Assistant Director File: 7-8-2 Joint Highway Research Project Project: 0-36-62B

Attached is a technical paper entitled, "Hydraulics of River Flow Under Arch Bridges -- A Progress Report, " which has treat breaged by Messrs. H. J. Owen, A. Socky, S. T. Husain, and Prof. J. W. Delleur. This paper was presented at the Joint Highway Resourch Project Sassion of the A5th Annual Purdue Road School.

The paper discusses the problem under study and the progress through Hapruary, 1959. This project is one of the cooperative projects which we have with the State Highway Department and the Eureau of Public Roads.

The paper will be published in the Proceedings of the 45th Attitual Road School and is presented to the Board for the record

Respectfully submitted.

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H. L. Michael, Secretary

HLM: ac

Attachment

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Technical Paper

HYDRAULICS OF RIVER FLOW UNDER ARCH BRIDGES -A PROGRESS REPORT-

Ву

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Joint Highway Research Project Project No: C-36-62B File No: 9-3-2

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HYDRAULICS OF RIVER FLOW UNDER ARCH BRIDGES -A PROGRESS REPORT-

By: H. J. Owen, A. Sooky, S. T. Husain, Graduate Assistants, Joint Highway Rosearch Project, and J. W. Delleur, Associate Professor of Hydraulic Engineering, Purdue University.

Introduction:

of a particular bridge depends on many variables. A major consider ion, often necessary, is the length of the bridge. In certain circum to an and locations the length may be dictated by factors other than cost. It, we are necessary the designer must decide, "How far can the bridge approach and tend onto the flood plain?". The shorter and therefore less expensive the bridge, the less waterway area generally provided. Waterway area in milmportant problem.

Bridge approaches extending far onto the flood plain decrease waterway area and produce, during high water, a large constriction consing oncessive backwater and possible damage to the structure as well to the necessary flooding of upstream areas. In some cases the state may be held liable for damage to property caused by bridge backwater.

In the past, studies pertaining to backwater caused by convertetions have considered shapes of opening such as that provided by a straight deck bridge. The Eureau of Public Roads has prepared, in cooperation with the Colorado State University, a report, entitled "Computation of Roadwater Caused by Bridges." This report in particular considered openings such as are provided by a straight deck bridge.

Figure 1 is a definition eketch for a normal crossing of the type used in the Colorado experiments. The typical water surface profile is shown as a solid line in view A. The maximum backwater superelevation is designated as $h_1^\#$ and the depression of the water surface below the normal

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http://www.archive.org/details/hydraulicsofrive00owen
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downstroam of the constriction as h3. Vica C of the figure illustrates the type opening studied.

Another remearch project whose subject is both interesting of applicable was carried out by New South Walon University of Technology. This project was primarily concerned with sharp-edged esctangular openings which produced contraction ratios, m, from 0.40 to 0.95. An equation was presented which gives the discharge as a function of the width of opening hase over the opening including backwater, y_1 , and 0 a coefficient depending on m and the Froude No. This equation makes it possible to determine the backwater if the discharge of the stream is known. Conversely, if the character of the discharge through the opening may be calculated by a measurement of y_1 .

Very little has been done with stoked openings. In short, to the knowledge of the authors, no systematic study has been used of the hydraulics of flow under such bridges. The arch bridge is unique in that the weileble waterway area decreases as the depth increases.

Figure 2 is a picture taken of the 19th Street bridge over Engle Gradi during the fleed occurring in the summer of 1957 at Speedway, Indiana. The taring effect of the arch and the accumulated debrie is visible. A typical multispan arch bridge subjected to flood flows is shown in Figure 3.

This bridge is the Wayne Street Bridge over the Wabash River at Peru, Indiana.*

A project was initiated in the Hydraulies Laboratory at Purdue University to study this problem. It is sponsored by the State Highway Department of Indians in cooperation with the U. S. Bureau of Public Roads.

Fhotographs courtesy of the Indiana State Flood Control and Water Resources Commission, Indianapolis.

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Furgoes:

The purpose, then, of the project is to:

- 1. Study the backwater produced by arches and develop a method for their computation.
- Pavelop criterion for designing the proper clear span.
- Study the hydraulic characteristics of flow under arch bridges including:
 - a) single span bridges
 - b) multiple span bridges
 - sequés decedudes bas noir analyses (o
 - d) shape of arch intrades
 - e) discharge and slope of stream
 - f) width of bridge.

This paper reports on the first year's work on this project. During this time, a preliminary investigation was initiated. Its purpose was fined to help in the design of the testing flume, and second to help in the design of the experiments to be carried out in the flume. Simultaneously the design of the required testing facilities was done and the constituction of those facilities was started.

Secre f Preliminary Experiments:

A dimensional analysis was made to define the important parameters of the flow. Originally, ten variables were considered as possibly of importance. Through the analysis the important parameters were found to be the Fronds Number, the roughness, the contraction ratio and the normal depth of the flow. The dimensional analysis is presented in Appendix I.

For the purpose of the preliminary testing, a small variable slope



flume 6" wide and 12' long was built. Figure 4 shows the laborator, equipment used in the preliminary testing. To the right of the figure, the fore-bay is visible. The channel sides and bottom were constructed on Profits and carefully aligned by manus of adjusting screws. The slope of the flume was controlled by a jack at the lower and of the flume. And be a writed horizontally above the flume served as a treat for the mechanics also triced point gages used in obtaining the water surface management.

An idealized two-dimensional case was investigated by the state circular weight with dismeter along the bottom as shown in Figure 1. represent the arch constriction. Sections B and C of Figure 5 illust the two types of surface profiles obtained with sill and steep sleps a squation relating the depths upstream of the weight and the discharge as developed in two ways. The exact solution presents Q in terms of the integrals:

$$Q = Cd \frac{4}{15} \sqrt{28} b^{5/2} \left\{ 2 \left(1 - k^2 + k^4 \right) \left[E - E(\phi, k) \right] - \left(1 - k^2 \right) \left(2 - k^2 \right) \left[k - E(\phi, k) \right] - r^2 \sin \phi \cos \phi \Delta \phi \left(3 k^2 \sin^2 \phi - 1 - k^2 \right) \right\}$$
(2)

Where K and E are the complete elliptic integrals of the first a 1 word kind respectively, $F(\varphi_j, E_j)$ and $E(\varphi_j, K)$ are the incomplete elliptic integrals of the first and second kind, and

$$k = \sqrt{\frac{y_1 + r}{b}} < 1$$
, $\phi = \sin^{-1} \sqrt{\frac{r}{y_1 + r}}$
 $\Delta \phi = \sqrt{1 - k^2 \sin^2 \phi} = \sqrt{0.5}$

An approximate solution gives Q in terms of an infinite series of presect of the ratio depth to radius:

$$Q = cy_1^{3/2}bA$$
 where $c = cd \frac{17}{24} \sqrt{29}$

and
$$A = \left[1 - 0.1294 \left(\frac{y_i}{r}\right)^2 - 0.0177 \left(\frac{y_i}{r}\right)^4, \dots\right]$$
 (2)



The Terrivotions of these formulas are given in Appendixes II and III.

The results of the wear tests were put in graphical form by plotting the coefficient of discharge vs the Fronds Number with the contraction rathers the parameter. The contraction ratio is defined as the model of the velo dismeter b to the fluxe width B. This graph is shown in the upper left scener of Figure 6. The lower graph shows the relation of the France Number of the ratio of depth upstream of the weir to the normal depth.

dim not had case by using sond-circular arch bridge models of the same contraction ratios. The bridge models were made of lubite-at typical water corract profile observation is shown in Figure 7. In that case, the fit is walls were limed with copper wire much of 16 meshes per inch. This gave a Manning's roughness coefficient of approximately 0.025, which is topical of many canals and natural streams. The general test results of the extension of the theory to the three-dissocional case obtained in the case the lucite channel and are shown in Figures 5 and 9.

Figure 8 shows the results of the three-dimensional tests, using bridge odels of width L = 2h inches. The coefficients of discharge 3d and the ratio of the backwater depth y_1 to the normal depth y_2 are plotted y_3 the facility Number for several values of the contraction ratio π . The results—if the two- and three-dimensional tests are compared in Figure 9. It is a teresting to notice that for small Froude Numbers, say less than 0.5. If and the ratios $\frac{y_1}{y_0}$ are approximately the same for the two cases. For higher Froude Numbers, the three-dimensional tests exhibit smaller values of 3d and larger values of $\frac{y_1}{y_0}$.

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Experimental Equipment and Test Plans:

designed and is now under construction. The large flure will be 5 ft wide, 64 ft. Ing and capable of 2 ft. maximum water depth. The structure is supported on six serew jacks with proportional rates of rise according to their positions. They will be driven by a common motor. These jacks will presit rapid and accurate changes of slope. Provision has been under for withning the absent eventually to 6 ft. Measurements of the water surface will be made first adjustable, stainless steed, guide rails running the length of the flume. The point gage to be used will be an electric indicating gage reading to a tenth of a millmeter. The flume will be provided with a tailgate control and a discharge control. Measurements of the discharge will be made thick two venturit tubes.

The models tested will first be confined to single spans with to skew. Later tests will include other variables. Assuming a typical bridge cross scatton the scale ratio between model and protetype will probably to from 1:6 to 1:15 for single span bridges.

And constant of the fluxe. The fluxe, as shown in the cross section essentially consists of two I beams for longitudinal support, then were 6° channels and rigidly attached vertical members. The specing between the channel members is 2 ft. Adjustment bolts are provided for leverage and alignment of the inner channel. The inner channel will be 1° plate steel.

This concludes the work done to date. This is a progress report on the first year of a three-year program.

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APPENDIX I

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Dimensional Analysia3

The Fuckingham Theorem states that in a physical problem including a quantities in which there are a dimensions, the quantities may be arranged into (n-m) dimensionless parameters. Suppose a dependent variable X_2 depends on the independent variables X_2 : X_1 X_n

$$\mathbf{x}_{1} = \mathbf{f}(\mathbf{x}_{2}, \mathbf{x}_{3}, \dots, \mathbf{x}_{n})$$

or
$$g(X_1, X_2, \ldots, X_n) = 0$$

If T_1 , T_2 ste, represent dimensionless grouping of the quantities X_1 , X_2 , X_3 , ste, with m dimensions involved, equation (4) may be replaced by an equation of the form

The method of obtaining the T parameters is to select m of the K quantities, with different dimensions and that contain among them the m dimensions. These m quantities are used as basic variables together with one of the remaining n-m quantities for each T. For example let K_1 , K_2 , K_3 contain M, L and T, not use assarily in each one, but collectively.

The first parameter is
$$T_{i} = X_{i} X_{2} X_{3} X_{4}$$
The second one
$$T_{2} = X_{i} X_{2} X_{3} X_{5}$$
and so in until
$$T_{n-m} = X_{i} X_{2} X_{3} X_{5}$$

In these equations the exponents are determined so that each is dimensionless. The dimensions of the X quantities are substituted and the exponents M, L. T are set equal to zero respectively. There will be three equations with three unknowns for each T parameter, so that the exponents a, b, c can be determined and hence the T parameters.

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In the problem at hand, it is desired to determine the backward, superelevation caused by an arch bridge constriction. It is assumed that the standables which govern this backwarder superelevation may be grouped into two categories: those describing the flew of the stream and those describing the bridge constriction. With reference to Figure 5, the following variables are considered:

- a) For the stream flow
 - y, maximum water elevation upstresm of constriction.
 - yo: the normal depth of flow in the approach channel
 - Vo. the velocity of flow at normal depth in the approach channel
 - n , Manning's roughness coefficient of the approach channel
 - $oldsymbol{arphi}$, the kinematic viscosity
 - P , density of the fluid
 - g, the acceleration of gravity
 - Ah, the maximum drop of the water surface caused by the construction
- b) For the structure—in order to simplify the problem let us consider at first single span circular arches with center on the bottom of the stream. The shape of the structure is thus fully determined by the diameter of the arch. The amount of constriction will also depend upon the initial width of the stream and the stage of the stream. The variables involved are thus:
 - B, width of stresm at the bridge site
 - b , bridge opening at the spring line
 - yo, normal depth.

Hence, from the above list of variables involved in our experiment:

$$y_{ij} = f(y_0, V, B, b, n, p, p, g, \Delta h)$$
 (5)

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	,	

or

There are m = 10 variables.

As n=3 dimensions are involved, M, L, U, are the their dimensions selected, and there are m-n = 10-3=7 factors. The backs variables so less of are V, ρ , y_0 , which are helpful for complex situations. The same

peranahora are:

by equaling the T quantities into dimensions

To (LT) (ML3) (L) LT-2

For L:

For T:

For M:

Solving:

Miones:

Similarly:

For L:

$$a = 3b + c = 1$$

For T:

For M:

Solving:

$$a = al_0 b = al$$
 and $c = al$

	'n	
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Thus:
$$T_2 = \frac{A^2}{V \, V_0}$$

Also: $T_3 = (LT^{-1})^{\alpha} (ML^{-3})^{\beta} L^{\alpha} L^{\gamma} L^{\gamma} = 0$

For L: $a - 3b + c + 1/6 = 0$

For T: $-a = 0$

For M: $b = 0$

Solving: $c = -1/6$

Then: $T_3 = \frac{a}{2\sqrt{6}}$

Also: $T_4 = (LT^{-1})^{\alpha} (ML^{-3})^{\beta} L^{\alpha} L^{\gamma}$

For L: $a - 3b + c + 1 = 0$

For M: $-a = 0$

For T: $+b = 0$

Solving: $c = -1$

Since B_{ϕ} b, heave linear dimensions einilar to y_{γ} :

$$T_5 = \frac{b}{30}$$
, $T_5 = \frac{b}{30}$, $T_7 = \frac{B}{30}$.

The sel

Introducing the W factors, equation (6) may be replaced by

$$h_1 = \frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}$$

Investing the first two parameters

Whence:

or
$$N_{1} = N_{2} = N_{1} = N_{2} = N_{2} = N_{3} = N_{4} = N$$

ber $\frac{\sqrt{g}}{\sqrt{g}}$. It is well known that gravity forces are predominant in open



charmed flow thereas viscous forces play a secondary role. The Republics mutter may therefore be disregarded. Furthermore, assuming that the shape of the material threshops of the material threshops of the material threshops of the water surface upstream, the term $\frac{a_0}{a_0 h}$ is also disregarded. Combining the ratio: $\frac{a_0}{a_0}$ and $\frac{a_0}{a_0}$ into $\frac{b}{b}$ equation (8) may be replaced by:

$$\frac{A}{40} = \phi \left[\frac{1}{5} \frac{2}{3}, \frac{3}{5} \frac{3}{5}, \frac{b}{5} \right] \tag{9}$$

It is thus seen that the backwater superclevation is appealed to be a .. metica of the Frends Number, the roughness, the contraction ratio and the normal depth of the flow.

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Equating (11) and (15) and solving for the condicions of a large

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Experimental cultimate of GL for Freedo Winterest , to 1 and 1 and of the consentation ratio (m = 0 050, $\kappa \approx 0.67$), m = 0 $\times 0$

Figure 6. From Abserved waltron of the decide up Q canter of the Property of the action of the property of the action of the action of the same values of the contraction ratio.



APPENDIK III



Transformation to Elliptic Integrals:

The integral of equation (10) may be evaluated in terms of complete and incomplete integrals of the first and second kind^{5,6}. Term Appendix III equation (10), the theoretical discharge Q_t is obtained making the coefficient of discharge Qd equal to unity:

$$Q_{\pm} = 2\sqrt{23} \int_{0}^{4\pi} \sqrt{(y_{1} - h)(r^{2} - h^{2})} dh$$
 (19)

and
$$\sin^2 u = \frac{h + r}{h + r}$$
 (21)

Since sulv. + culv. a j

since d (snu) = snu dau

From (21) and making use of (20)

Also from (21) $n^2 - h^2 = 4r^2k^2 sn^2u \left(1 - k^2 sn^2u\right)$ $= 4r^2k^2 sn^2u dn^2u$ (25)

since dn2u + k2sn2u = 1

Substituting (23), (24) and (25) into (19) the expression for the theoretical discharge becomes

Charge becomes

$$Q_{2} = 2\sqrt{2g} \int_{u_{1}}^{u_{2}} \left[2rk^{2}cn^{2}u \cdot 4r^{2}k^{2}sn^{2}udn^{2}u \right] 4rk^{2}snudnucnudu$$

(26)

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(3)

The lower limit u, is obtained from (21) as follows:

or snu, =
$$\frac{0+n}{41+n}$$

or snu, = $\sin \phi = \sqrt{\frac{n}{41+n}}$
 $\phi = am \omega_1$

and finally
$$w_i = F(\phi, k) = \int_0^{\phi} \frac{d\phi}{\sqrt{1 - k^2 \epsilon^2 k^2 \epsilon^2 k^2 \epsilon^2}}$$
, is (27)

 $F(\phi_1 k)$ is the incomplete elliptic juty grat of the p=1. The upper limit u_2 , is obtained from (21) as follows:

or

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where K is the complete elliptic integral of the first kind 23

The empression for the theoretical diachs (20 - 26) becomes

$$Q_{\pm} = 32\sqrt{g} V^{5/2} L^4 \int_{-1}^{1} C n^3 u \, dn^2 u \, sn^3 \, u \, dn$$

Upon performing the integration, and introducing the diameter
$$b = 2r$$

$$Q_{2} = \frac{2}{15}\sqrt{2}q_{1}b^{3/2}\left[2\left(1-k^{2}+16^{4}\right)\left[2-2\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left[12-7+k^{2}\right]\right]$$

$$= \frac{1}{15}\sqrt{2}q_{1}b^{3/2}\left[2\left(1-k^{2}+16^{4}\right)\left[2-2\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{2}\right)\left(1-k^{$$

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$$E = \int_{0}^{\pi/2} \sqrt{1 - \mu^{2} \sin^{2}\phi} d\phi \tag{31}$$

which is the complete elliptic integral on the second kinc, and

$$E(\phi_1 k) = \int_0^{\phi} \sqrt{1 - k^2 \sin^2 \phi} dy$$
 (32)

which is the incomplete elliptic integral of the second kind, and

$$\phi = \sin^{-1}\sqrt{\frac{c}{y_1+n}} \tag{33}$$

and

$$\Delta \phi = \sqrt{1 - k^2 \sin^2 \phi} = \sqrt{0.5} \tag{34}$$



and finally

Equation (30) yields the theoretical discharge for the flow through a semicircular constriction of dismeter b=2r and where the maximum death upstream of the constriction is y_1 . The quantities K, E, F (φ , k), E (φ , k) may be obtained from tables. Equation (30) is compated similar to that for the flow through circular weights obtained by F. C. Stevens F which is

$$Q_{\pm} = \frac{4}{15} \sqrt{29}, \quad \int_{-10}^{1/2} \left(2(1-k^{5}+k^{4}) E - (2-k^{2})(1-k^{4}) E \right)$$
 (35)

where $k^2 = H/D$, H being the head over the invert, and D is the district of the circular weir. Stevens also gives an infinite series approximation to equation (36) which is similar to equation (11):

$$Q_{6} = 2\sqrt{2}q. D^{5/2}T\left(\frac{1}{6}z^{2} - \frac{1}{2}z^{3} - \frac{5}{10.37}z^{4} + \dots\right)$$
(3.1)

where z = H/D.



SYMBOLS

	A	An infinite series of powers of the ratio depth to radius.
	dA	Elementary area
	B	Width of arch at spring line, or with
	ો	Diameter of the arch or voir.
	C	Coefficient dap using on the coefficien of discharge, equation (25).
	Cá	Coefficient of cischarge
	正	Complete elaiptic integral of the second kind, equation (31).
	E (\$, k)	Incomplete illiptic integral of the second kind, equation (32)
	<u>[4]</u> 37	Froude Number of the flow, equation (16)
	? (ф, k)	Incomplete illiptic integral of the second kind, equation (32).
	E	Acceleration of gravity.
	Î.	Hydraulic hesd.
	h hj	Maximum backweter superelevation.
	-}}-	
L	h ₁	Maximum backwater superplevetion. Depression of the water surface below the
	h ₁	Maximum backweter superclevetion. Depression of the water surface below the normal dometream of the constriction. Drop in water surface caused by the con-
Δ	h ₁ h ₂	Maximum backweter superelevation. Depression of the water surface below the normal downstream of the constriction. Drop in water surface caused by the constriction. Complete alliptic integral of the first
8.2	h ₁ h ₃ h ₃	Maximum backwater superplevation. Depression of the water surface below the normal domestream of the constriction. Drop in water surface caused by the constriction. Complete elliptic integral of the first kind, equation (28).
L	h ₁ h ₂ h ₃ k	Maximum backweter superclevation. Depression of the water surface below the normal downstream of the constriction. Drop in water surface caused by the constriction. Complete elliptic integral of the first kind, equation (28). Medulus of elliptic integrals, equation (20).
L	h ₁ h ₂ h ₃ h ₃ k	Meximum backweter superclavation. Depression of the water surface below the negmal downstream of the constriction. Drop in water surface caused by the constriction. Complete elliptic integral of the first kind, equation (28). Medulus of elliptic integrals, equation (20). Width of bridge, measured along waterway axis.
	h ₁ h ₃	Maximum backweter superelevation. Depression of the water surface below the negmal downstream of the constriction. Drop in water surface caused by the constriction. Complete elliptic integral of the first kind, equation (28). Medulus of elliptic integrals, equation (20). Width of bridge, measured along waterway axis. Contraction ratio = b/B.



Q	Actual dischargo.
O _t	Theoretical discharge.
χ.,	Radius of the arch.
sn u, en u, dn u	Jacobi elliptic functions.
$V_{\mathbf{O}}$	Velocity at normal depth in approach channel.
Y.	Normal depth in approach channel.
71	Maximum water depth upstress of constriction.
\checkmark	Kinematic viscosity.
P	Density of the fluid.
ø	Amplitude in elliptic functions, $\phi = sm u$, equation (33).



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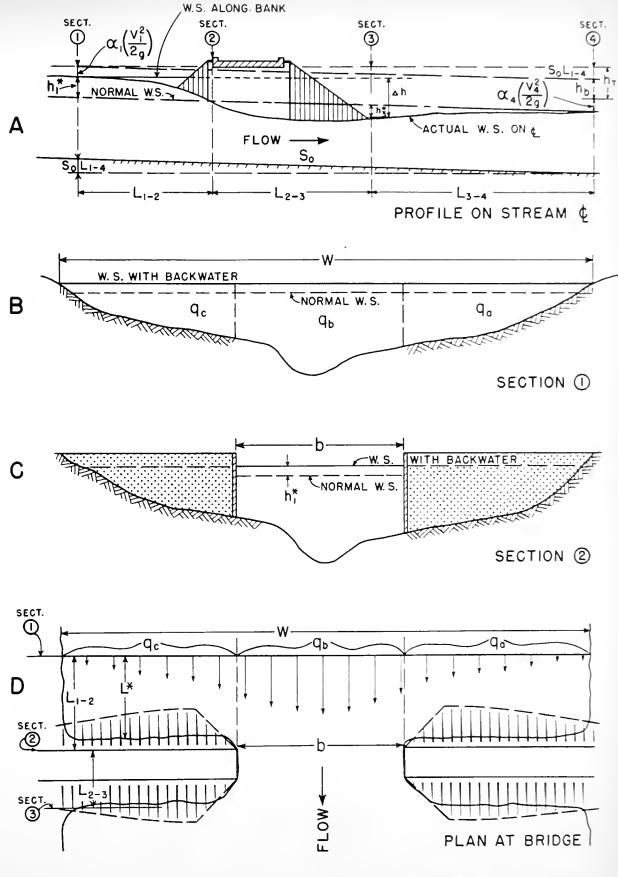


FIG. I

NORMAL CROSSING WINGWALL ABUTMENTS



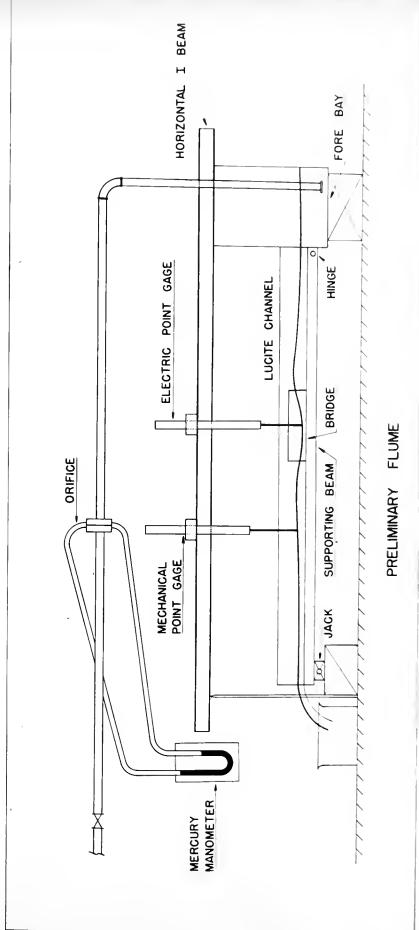


FIG. 2



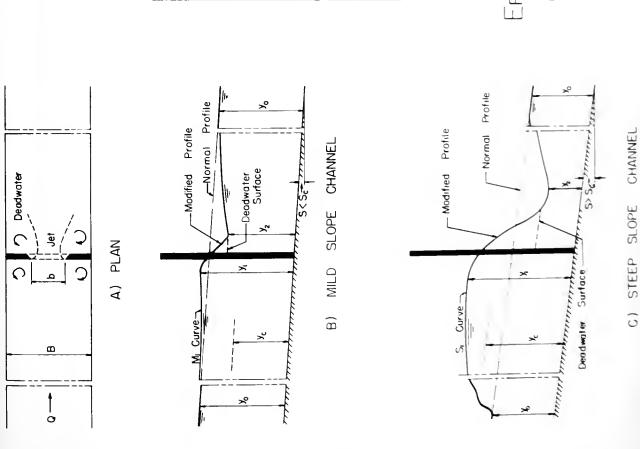
FIG. 3



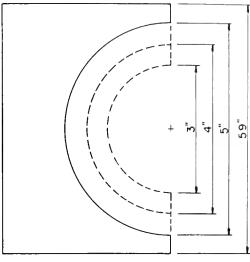


F1G. 4





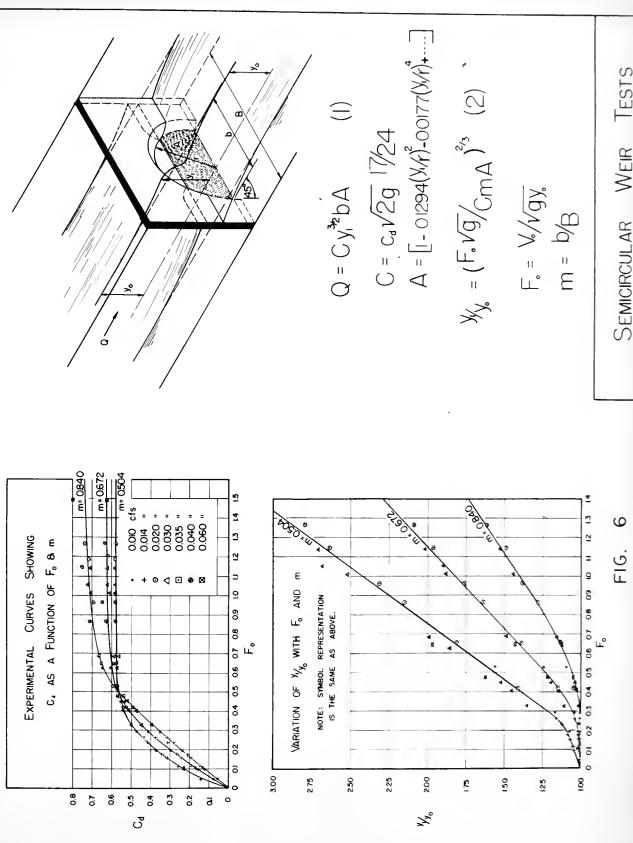
D.) WEIR PLATES



EFFECT OF CHANNEL CONSTRICTION ON WATER SURFACE PROFILE

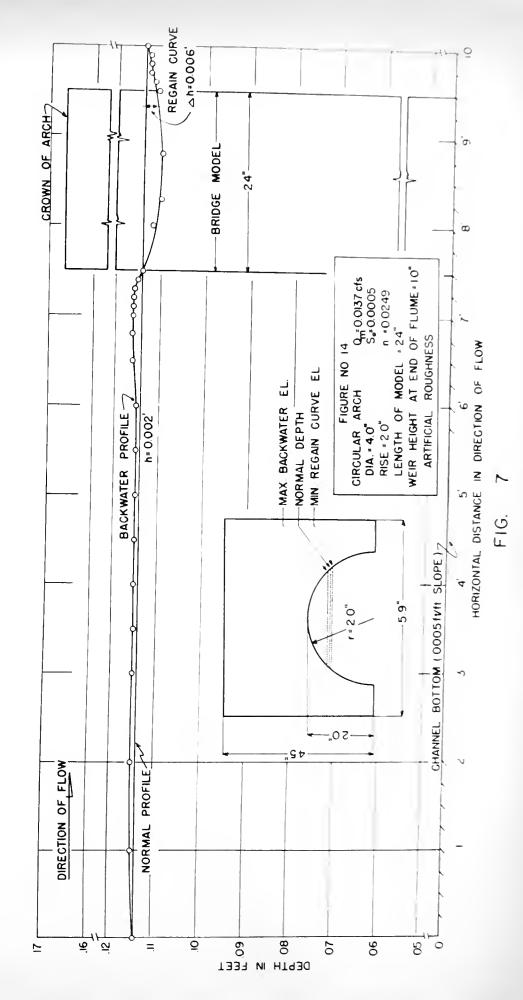
F1G. 5



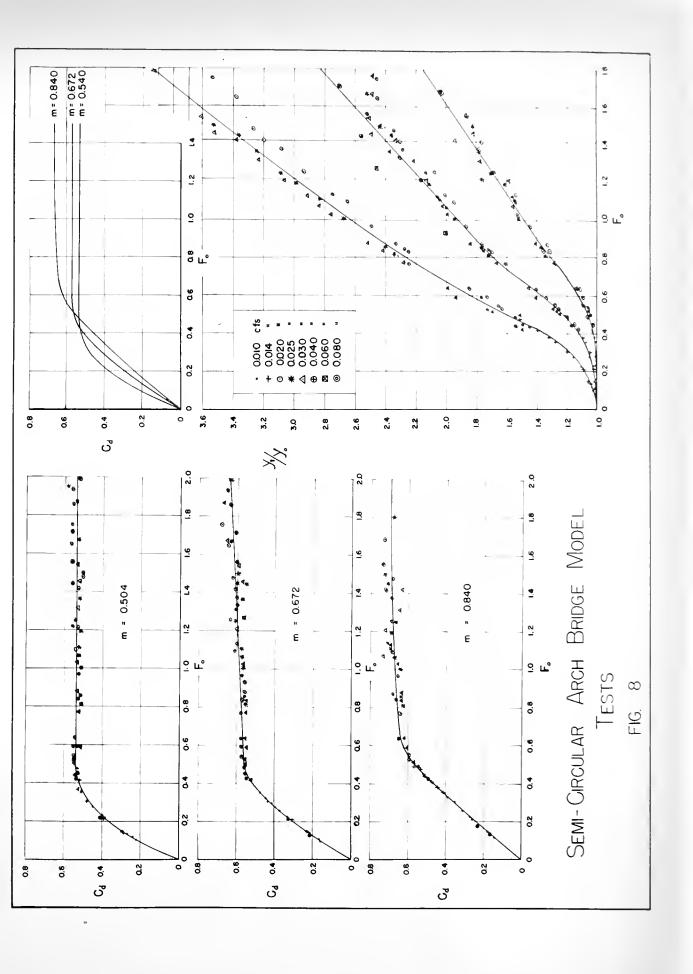


SEMICIRCULAR WEIR TESTS

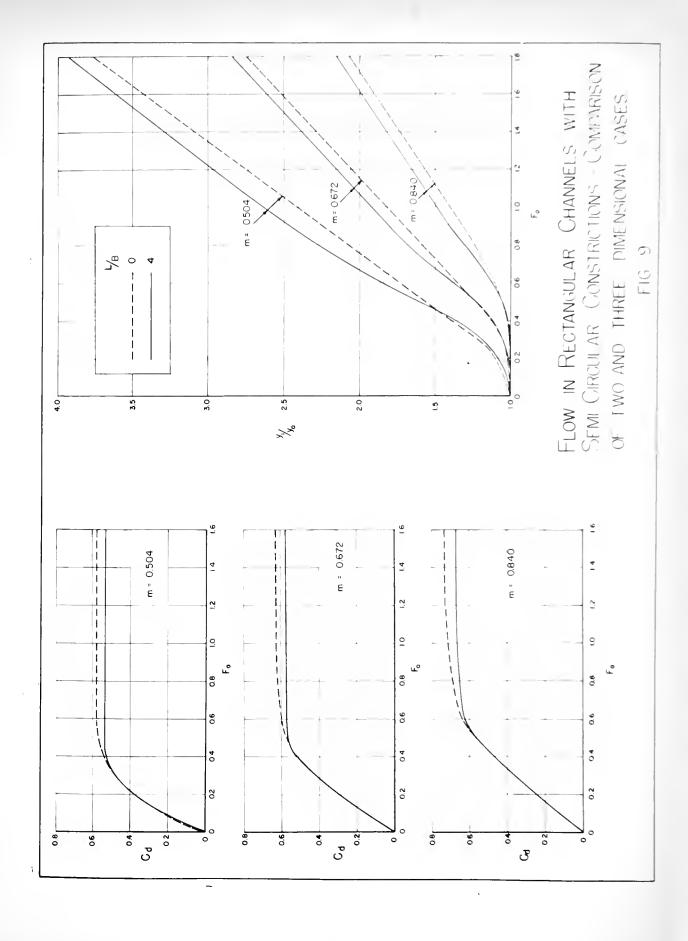




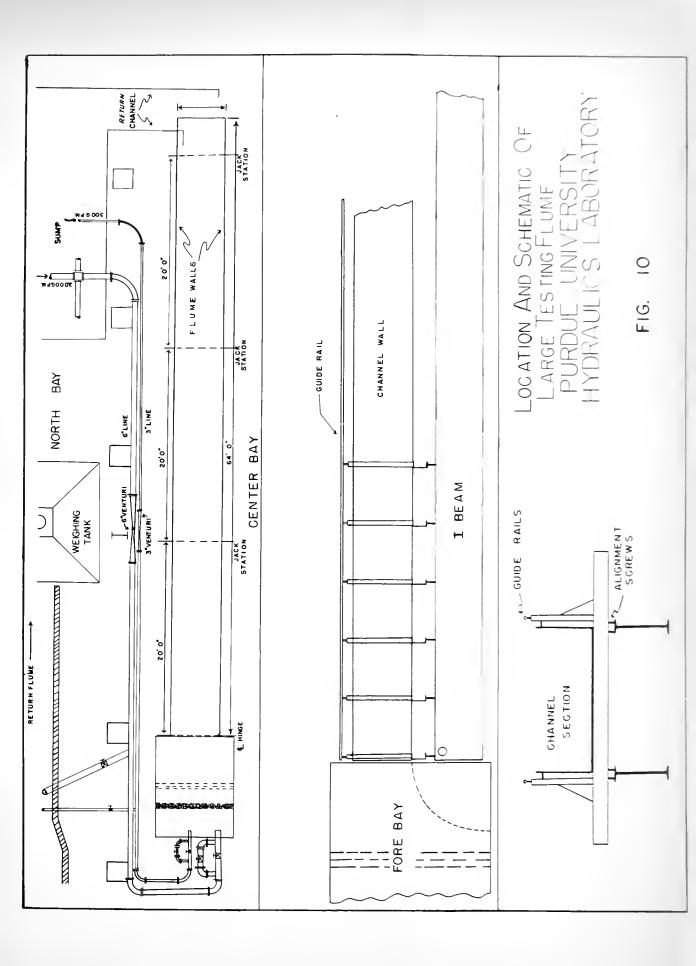


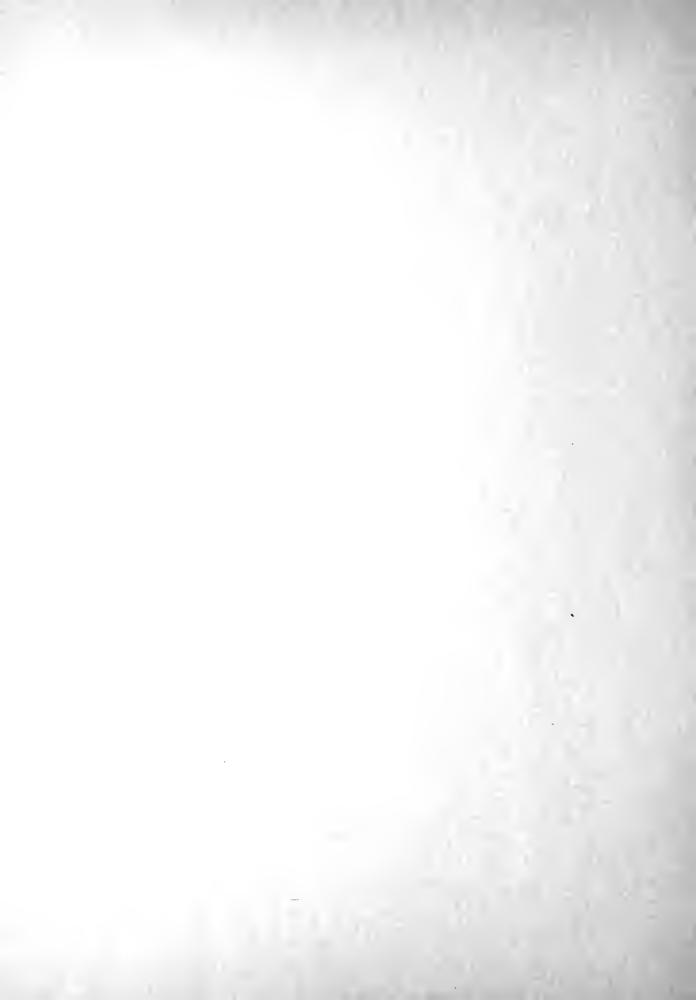


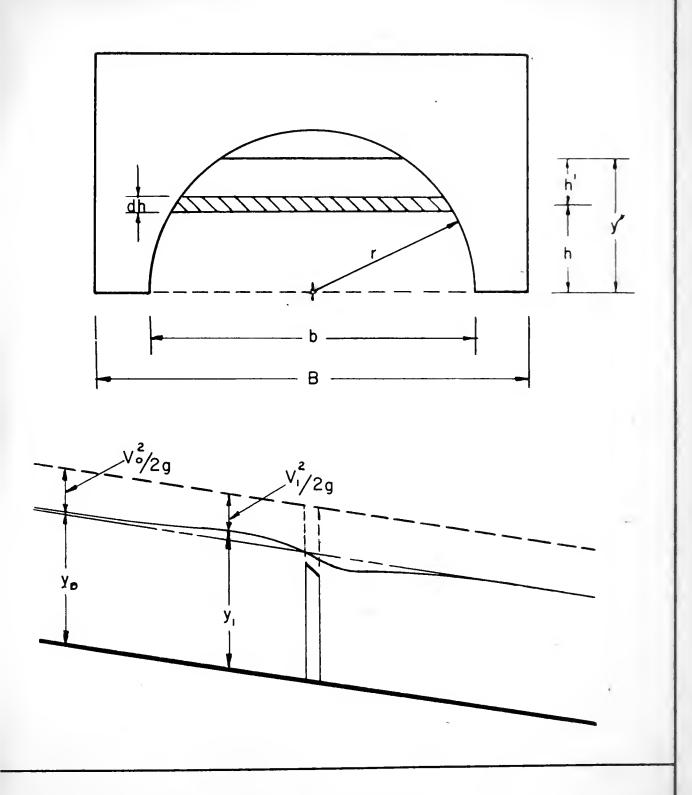












DEFINITION SKETCH

FIG. 11

